General Relativity

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Exercise 1: Relations for the metric tensor

In the lecture we postulated the gauge transformation of the metric tensor $g_{\mu\nu}$,

$$\delta_{\xi}g_{\mu\nu} = \xi^{\rho}\partial_{\rho}g_{\mu\nu} + \partial_{\mu}\xi^{\rho}g_{\rho\nu} + \partial_{\nu}\xi^{\rho}g_{\mu\rho} .$$
(1)

- 1.1) Compute the gauge transformation of the inverse metric $g^{\mu\nu}$.
- 1.2) Prove that the determinant $|g| \equiv \det g_{\mu\nu}$ and the inverse metric $g^{\mu\nu}$ can be written as

$$egin{array}{rcl} |g| &=& rac{1}{4!} \, arepsilon^{\mu
u
ho\sigma} \, arepsilon^{lphaeta\gamma\delta} g_{\mulpha} \, g_{
ueta} \, g_{
ho\gamma} \, g_{\sigma\delta} \; , \ g^{\mu
u} &=& rac{1}{3!} \, rac{1}{|g|} \, arepsilon^{\mu
ho\sigma\kappa} \, arepsilon^{
ulphaeta\gamma} \, g_{
holpha} \, g_{\sigmaeta} \, g_{\sigma\kappa\gamma} \; , \end{array}$$

where $\varepsilon^{\mu\nu\rho\sigma}$ is the totally antisymmetric Levi-Civita symbol with $\varepsilon^{0123} = +1$.

1.3) Determine $\delta \sqrt{-g}$, where $g \equiv |g|$, under arbitrary variations $\delta g_{\mu\nu}$.

Exercise 2: Lie derivatives

Lie derivatives are defined as

$$\mathcal{L}_{\xi}V^{\mu} = \xi^{\nu}\partial_{\nu}V^{\mu} - \partial_{\nu}\xi^{\mu}V^{\nu} , \qquad \mathcal{L}_{\xi}W_{\mu} = \xi^{\nu}\partial_{\nu}W_{\mu} + \partial_{\mu}\xi^{\nu}W_{\nu} ,$$

and analogously for objects with an arbitrary number of upper and lower indices.

2.1) Prove that the Lie derivative operators \mathcal{L}_{ξ} satisfy the algebra

$$[\mathcal{L}_{\xi_1}, \mathcal{L}_{\xi_2}] = \mathcal{L}_{[\xi_1, \xi_2]},$$

where [,] on the left-hand side denotes the commutator. In the process, determine the Lie bracket $[\xi_1, \xi_2]$ defined by the right-hand side.

2.2) We defined the *non-covariant variation* $\Delta_{\xi} \equiv \delta_{\xi} - \mathcal{L}_{\xi}$, the difference of the actual gauge transformation and the covariant one given by the Lie derivative. Compute

$$\Delta_{\xi}\Gamma^{\rho}_{\mu\nu}$$
 and $\Delta_{\xi}(\partial_{\mu}\Gamma^{\rho}_{\nu\sigma})$,

where $\Gamma^{\rho}_{\mu\nu}$ are the Christoffel symbols for $g_{\mu\nu}$, and express the result, to the extent possible, with matrix notation for $\Gamma_{\mu} \equiv (\Gamma_{\mu})^{\rho}{}_{\nu} \equiv \Gamma^{\rho}_{\mu\nu}$ and $\sigma \equiv \sigma^{\mu}{}_{\nu} \equiv \partial_{\nu}\xi^{\mu}$.

Exercise 3: Expansion of non-linear gravity around flat space

In the lecture it was claimed that the following action describes non-linear gravity:

$$S = \int d^4x \sqrt{-g} \left(\frac{1}{4} g^{\mu\nu} \partial_{\mu} g^{\rho\sigma} \partial_{\nu} g_{\rho\sigma} - \frac{1}{2} g^{\mu\nu} \partial_{\mu} g^{\rho\sigma} \partial_{\rho} g_{\sigma\nu} \right. \\ \left. + \partial_{\mu} g^{\mu\nu} \frac{1}{\sqrt{-g}} \partial_{\nu} \sqrt{-g} + g^{\mu\nu} \left(\frac{1}{\sqrt{-g}} \partial_{\mu} \sqrt{-g} \right) \left(\frac{1}{\sqrt{-g}} \partial_{\nu} \sqrt{-g} \right) \right)$$

(3.1) Expand the above action around Minkowski space to second order in h, writing

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu} ,$$

and show that it reproduces the Fierz-Pauli action. Hint: Use the result of 1.3) in order to expand $\sqrt{-g}$.

3.2) Take the above expansion of $g_{\mu\nu}$ to be exact (i.e. no higher order terms of h are included on the right-hand side). Determine the complete gauge transformation $\delta_{\xi}h_{\mu\nu}$ for the fluctuation field from eq. (1) above.